

FIG. 2.4. Geometrical interpretation of the relation between the Gibbs and Helmholtz potentials at a fixed temperature  $T > T_c$ . This construction, adapted from Morse (1969) and Griffiths (unpublished lecture notes), is based on the identities  $A = G - PV$  and  $V = (\partial G / \partial P)_T$ . The vertical distance between the dashed lines in (a) and (b) is the product  $PV$ , and on subtracting this from  $G$  we obtain  $A$ . Note that the construction indicated is unambiguous because  $G(T, P)$  is a concave function of  $P$  (for all  $P$ ) and  $A(T, V)$  is a convex function of  $V$  (for all  $V$ ). Also shown are the volume as a function of pressure (obtained from the pressure derivative of  $G(T, P)$ ) and the pressure as a function of volume (obtained from the volume derivative of  $A(T, V)$ ).

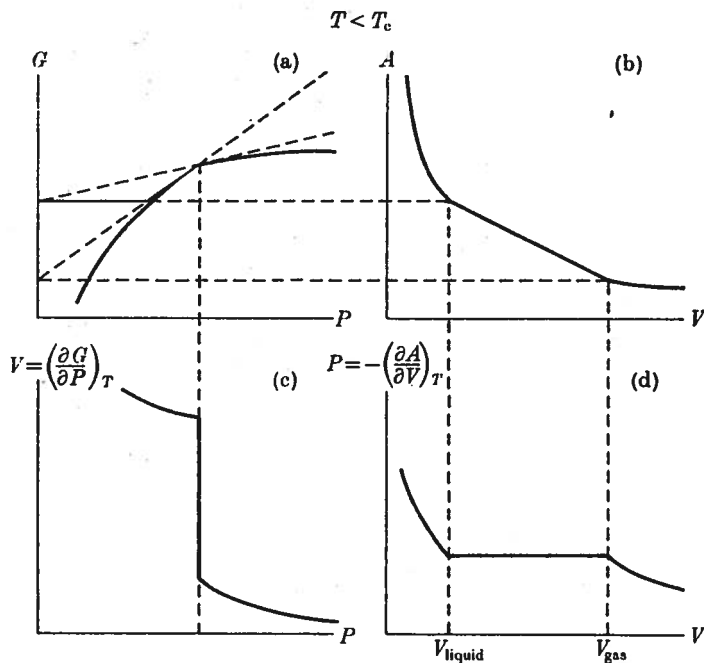


FIG. 2.5. The Gibbs and Helmholtz potentials for a fixed temperature  $T < T_c$ , where  $T_c$  is the transition temperature for a first-order phase transition. Also shown are the volume as a function of pressure and the pressure as a function of volume. Note that the value of the pressure at which the volume discontinuity occurs is not the critical pressure  $P_c$  but rather the value of the saturated vapour pressure  $P_{\text{sat}}(T)$  given by the vapour pressure curve of Fig. 1.1a. Thus the system is seen to undergo a *first-order* transition on crossing the vapour pressure curve at constant temperature  $T$  ( $T < T_c$ ).

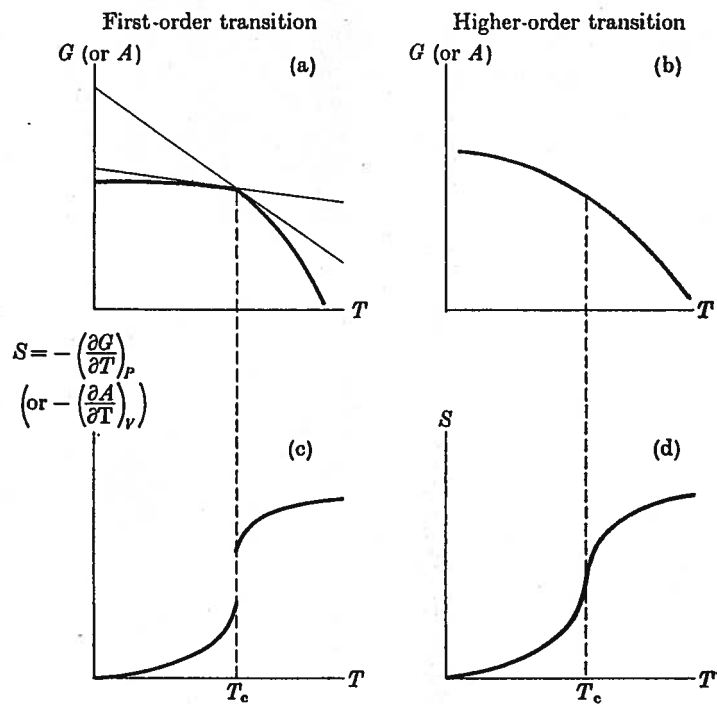


FIG. 2.6. (a) Temperature dependence of either the Gibbs potential at fixed pressure or the Helmholtz potential at fixed volume. The system shown undergoes a phase transition at  $T = T_c$ , accompanied by a latent heat (entropy discontinuity). (b) Same as (a) except that the phase transition has no latent heat (entropy discontinuity). (c)-(d) show the entropy obtained from the temperature derivative of  $A$  or  $G$ . For this particular example,  $C \sim dS/dT$  appears to diverge at  $T_c$  for both the first-order and the higher-order transition.

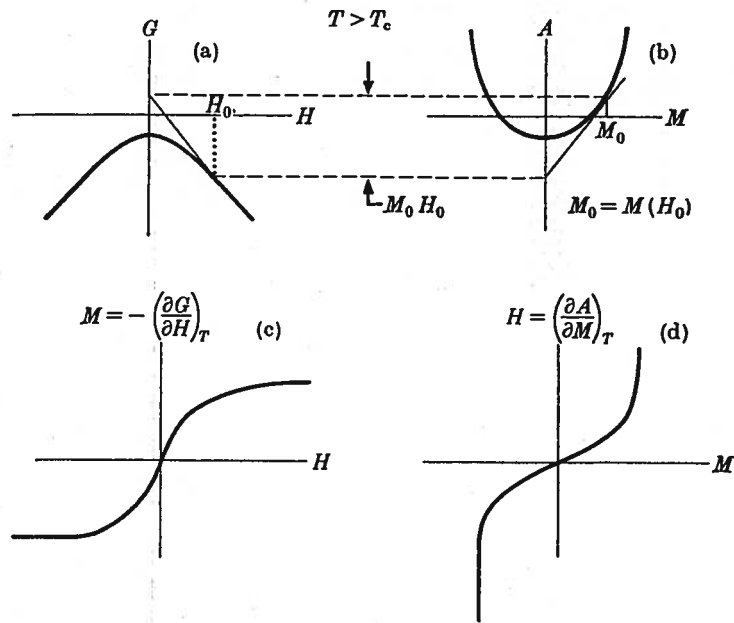


FIG. 2.8. Geometrical relation between  $G(T, H)$  and  $A(T, M)$  for a fixed temperature  $T > T_c$ . Also shown are the magnetization as a function of field (obtained from  $M = -(\partial G/\partial H)_T$ ) and the field as a function of magnetization (obtained from  $H = (\partial A/\partial M)_T$ ).

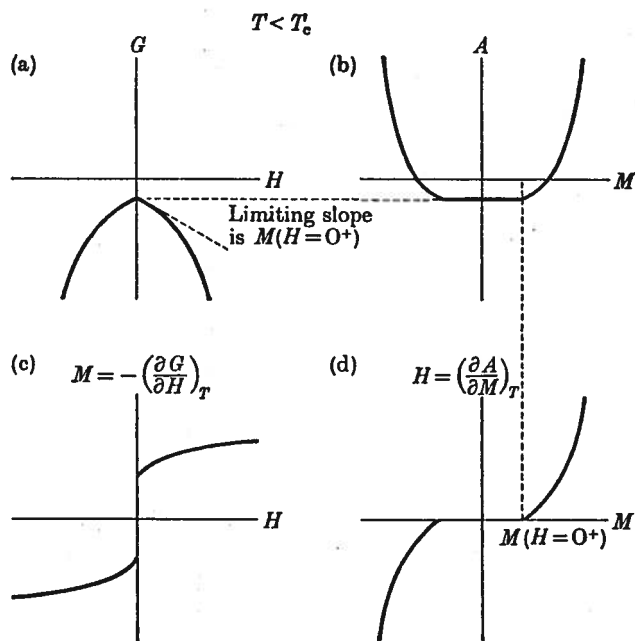


FIG. 2.9. Relation between  $G(T, H)$  and  $A(T, M)$  for a fixed temperature  $T < T_c$ . Also shown are the  $M$ - $H$  isotherms (obtained from  $G(T, H)$ ) and the  $H$ - $M$  isotherms (obtained from  $A(T, M)$ ).